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# Isotopically Substituted Bis(ethylenedithio)tetrathiafulvalene at 15 K by Time-of-Flight Neutron Diffraction 

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#### Abstract

The title compound, $\mathrm{C}_{10} \mathrm{D}_{8} \mathrm{~S}_{8}$, (ET), has been synthesized whereby all eight S atoms, the four ethylene C atoms and all eight H atoms have been isotopically substituted with ${ }^{34} \mathrm{~S},{ }^{13} \mathrm{C}$ and ${ }^{2} \mathrm{H}$, respectively. Analysis of the neutron scattering lengths derived from singlecrystal time-of-flight neutron diffraction data obtained at $T=15 \mathrm{~K}$ confirms the isotopic distribution. The title phase crystallizes in the space group $P 2_{1} / n$ with bond lengths and angles consistent with neutral ET.


## Comment

The title compound, (I), (ET), participates as the electron-donor molecule in a large number of organic superconducting salts (Williams et al., 1991). In order to investigate the mechanism for the electron pairing in this class of superconductors, a variety of isotopically substituted ET molecules with different distributions of isotopes throughout the molecule have been synthesized (Kini et al., 1996). Salts of these organic donor molecules have been studied to determine the isotopic effect on $T_{c}$ in the superconducting crystals. In the crystal used in this neutron diffraction experiment, the isotopic substitution of the ET molecule involved the replacement of the four C atoms in the terminal 1,2 -ethanediyl groups with ${ }^{13} \mathrm{C}$, all eight S atoms with ${ }^{34} \mathrm{~S}$ and all eight H atoms with deuterium, ${ }^{2} \mathrm{H}$. The substitution increases the mass of the ET molecule, hereafter denoted by ${ }^{2} \mathrm{H}_{8}{ }^{13} \mathrm{C}_{4}{ }^{34} \mathrm{~S}_{8}$-ET, by 28 a.m.u., which is about $7 \%$ of the original mass. From the study by Kini et al. (1996), it was shown that the C and S substitution has a normal BCS-like mass ( $M$ ) effect ( $T_{c}$ proportional to $M^{-\alpha}$ ),
whereas deuterium substitution produces an inverse isotopic effect ( $T_{c}$ increases). We present here the neutron structure of ${ }^{2} \mathrm{H}_{8}{ }^{13} \mathrm{C}_{4}{ }^{34} \mathrm{~S}_{8}$-ET at 15 K in order to confirm the isotopic distribution and to provide high-precision low-temperature structural parameters.

(I)

The only previous structure of neutral ET is the roomtemperature single-crystal X-ray structure in space group $P 2_{1} / c$ by Kobayashi, Kobayashi, Yukiyoshi, Saito \& Inokuchi (1986). We had transformed the atomic coordinates from $P 2_{1} / c$ to $P 2_{1} / n$ in our previous comparison of the ET structure with the isomorphous $\mathrm{Au}(\mathrm{DDDT})_{2}$ (DDDT $=5,6$-dihydro-1,4-dithiin-2,3-dithiolate) crystal structure (Schultz et al., 1987). The transformation of the atomic coordinates from the $P 2_{1} / c$ unit cell in Kobayashi et al. (1986) to the $P 2_{1} / n$ unit cell in this paper is $(x-z,-y,-z)+(1 / 2,1 / 2,0)$.

The $P 2_{1} / n$ room-temperature unit cell of neutral ET reported by Kobayashi et al. (1986) is $a=6.614$ (11), $b=13.985$ (2), $c=15.721$ (3) $\AA, \beta=93.81$ (2) ${ }^{\circ}$ and $V$ $=1451(5) \AA^{3}$. The low-temperature data of isotopically substituted ET gave a slightly smaller volume, as expected, of 1416 (1) $\AA^{3}$ and $a=6.661$ (1), $b=13.652$ (2), $c=15.613$ (2) $\mathrm{A}, \beta=94.281$ (2) ${ }^{\circ}$. The bond lengths of the low-temperature isotopically substituted ET appear to be slightly longer than those obtained from room-temperature data. Shorter bond distances of the room-temperature neutral ET are probably attributable to anisotropic thermal motion. Isotopically substituted ET bond distances at 15 K are $\mathrm{C} 1-\mathrm{C} 2=1.359$ (2), $\mathrm{C} 5-\mathrm{C} 6=1.360(2), \mathrm{S} 2-\mathrm{C} 4=1.771$ (3) and S4-C6 $=1.769$ (3) A. Kobayashi et al. (1986) reported 1.319, $1.328,1.757$ and $1.751 \pm(0.007-0.012) \AA$, respectively.

The ET molecules pack as face-to-face dimeric pairs. Many of the S-S contacts between different molecules have distances close to the sum of the van der Waals radii $(3.70 \AA)$. The $\mathrm{S} \cdots \mathrm{S}$ contacts are listed in Table 2.
The crystal for this data collection was obtained from a batch of crystals which underwent mass spectrometry to verify isotopic substitution. In order to confirm the degree of isotopic substitution of this crystal, scattering lengths for each of the atoms were refined. Scattering lengths in units of $10^{-12} \mathrm{~cm}$ from Sears (1986) are: $b\left({ }^{12} \mathrm{C}\right)=0.66511(16), b\left({ }^{13} \mathrm{C}\right)=0.619(9), b\left({ }^{32} \mathbf{S}\right)=$ $0.2804(2), b\left({ }^{34} \mathrm{~S}\right)=0.348(3), b(\mathrm{H})=-0.37406(11)$, $b(\mathrm{D})=0.6671$ (4). The average refined scattering length in this structure of the ethanediyl ${ }^{13} \mathrm{C}$ atoms is 0.656 (3) $\times 10^{-12} \mathrm{~cm}$, which appears to be closer in magnitude to ${ }^{12} \mathrm{C}$ but is also within three standard deviations of the ${ }^{13} \mathrm{C}$ scattering length. Thus, the refined ethanediyl
scattering lengths do not statistically verify the ${ }^{13} \mathrm{C}$ substitution and do not permit an assignment of the C isotopes. The deuterium substitution is $96(2) \%$ and is obtained from an average scattering length of 0.621 (2) $\times 10^{-12} \mathrm{~cm}$. Sulfur sites are completely substituted with ${ }^{34} \mathrm{~S}$. The average scattering length is $0.354(2) \times$ $10^{-12} \mathrm{~cm}$, indicative of $108(5) \%$ substitution of ${ }^{34} \mathrm{~S}$.


Fig. 1. Molecular structure of isotopically substituted ET with displacement ellipsoids drawn at the $75 \%$ probability level.

## Experimental

The title compound was prepared according to Kini et al. (1996).

## Crystal data

$\mathrm{C}_{10} \mathrm{D}_{8} \mathrm{~S}_{8}$
$M_{r}=412.6$
Monoclinic
$P 2_{1} / n$
$a=6.661(1) \AA$
$b=13.652(2) \AA$
$c=15.613$ (2) $\AA$
$\beta=94.281(2)^{\circ}$
$V=1416(1) \AA^{3}$
$Z=4$
$D_{x}=1.93 \mathrm{Mg} \mathrm{m}^{-3}$
$D_{m}$ not measured

## Data collection

IPNS single-crystal diffractometer
Time-of-flight data
, Absorption correction: spherical
$T_{\text {min }}=0.9960, T_{\text {max }}=$ 0.9961

Refinement
Refinement on $F$
$R=0.062$
$w R=0.058$
$S=1.2$
$4571[F<3 \sigma(F), d>$ $\left.0.7 \AA, F_{o}^{2}<2 F_{c}^{2}\right]$ reflections
274 parameters
All D atoms refined $w=4 F^{2} /\left[\sigma^{2}\left(F^{2}\right)\right.$
$\left.+\left(0.04 F^{2}\right)^{2}\right]$

Pulsed neutron radiation
$\lambda=0.7-4.2 \AA$
Cell parameters from 1941 reflections
$\mu=0.0355$ $+0.00052(\lambda) \mathrm{mm}^{-1}$
$T=15 \mathrm{~K}$
Rectangular
$3.5 \times 1.5 \times 1.0 \mathrm{~mm}$ Red


19444 measured reflections 9580 independent reflections 5405 reflections with
$F>3 \sigma(F)$
$h=-12 \rightarrow 13$
$k=-19 \rightarrow 9$
$l=-11 \rightarrow 22$
$(\Delta / \sigma)_{\max }=0.02$
$\Delta \rho_{\text {max }}$ equivalent to $1 \%$ of an $S$ atom
$\Delta \rho_{\text {min }}$ equivalent to $1 \%$ of an $S$ atom
Extinction correction: Becker \& Coppens (1974) type I
Extinction coefficient:

$$
0.91(2) \times 10^{4}
$$

Scattering factors from Sears (1986)

Table 1. Fractional atomic coordinates and equivalent isotropic displacement parameters $\left(\AA^{2}\right)$

| $U_{\text {eq }}=(1 / 3) \sum_{i} \sum_{j} U^{i j} a_{i}^{*} a_{j}^{*} \mathbf{a}_{i} \cdot \mathbf{a}_{j}$. |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $x$ | $y$ | $z$ | $U_{\text {eq }}$ |
| S1 | -0.2580 (4) | 0.1167 (2) | -0.0207 (2) | 0.0048 (7) |
| S2 | 0.1719 (4) | 0.1417 (2) | 0.0421 (2) | 0.0052 (7) |
| S3 | -0.3416 (4) | -0.0046 (2) | 0.1534 (2) | 0.0052 (7) |
| S4 | 0.0890 (4) | 0.0116 (2) | 0.2124 (2) | 0.0046 (7) |
| S5 | -0.2434 (4) | 0.2769 (2) | -0.1531 (2) | 0.0062 (7) |
| S6 | 0.2332 (4) | 0.3039 (2) | -0.0878 (2) | 0.0052 (7) |
| S7 | -0.4796 (4) | -0.0673 (2) | 0.3203 (2) | 0.0067 (8) |
| S8 | 0.0366 (4) | -0.0507 (2) | 0.3929 (2) | 0.0057 (7) |
| C1 | -0.0686 (2) | 0.0966 (1) | 0.0628 (1) | 0.0061 (4) |
| C2 | -0.1031 (2) | 0.0439 (1) | 0.1342 (1) | 0.0062 (4) |
| C3 | -0.1212 (2) | 0.2096 (1) | -0.0696 (1) | 0.0060 (4) |
| C4 | 0.0737 (2) | 0.2219 (1) | -0.0398 (1) | 0.0065 (4) |
| C5 | -0.2757 (2) | -0.0270 (1) | 0.2642 (1) | 0.0055 (4) |
| C6 | -0.0774 (2) | -0.0198 (1) | 0.2911 (1) | 0.0065 (4) |
| C7 | -0.0499 (3) | 0.2891 (1) | -0.2315 (1) | 0.0086 (5) |
| C8 | 0.1663 (3) | 0.2662 (1) | -0.1989 (1) | 0.0070 (5) |
| C9 | -0.3618 (3) | -0.0783 (1) | 0.4285 (1) | 0.0073 (5) |
| C10 | -0.1570 (2) | -0.1288 (1) | 0.4319 (1) | 0.0080 (5) |
| D7A | -0.0647 (3) | 0.3652 (1) | -0.2526 (1) | 0.0207 (6) |
| D7B | -0.0956 (3) | 0.2409 (2) | -0.2856 (1) | 0.0216 (6) |
| D8A | 0.1977 (3) | 0.1875 (1) | -0.2011 (1) | 0.0185 (6) |
| D8B | 0.2719 (3) | 0.3048 (1) | -0.2381 (1) | 0.0181 (6) |
| D9A | -0.4695 (3) | -0.1218 (1) | 0.4626 (1) | 0.0174 (6) |
| D9B | -0.3478 (3) | -0.0049 (1) | 0.4575 (1) | 0.0168 (6) |
| D10A | -0.1046 (3) | -0.1431 (1) | 0.4990 (1) | 0.0171 (6) |
| D10B | -0.1640 (3) | -0.1984 (1) | 0.3958 (1) | 0.0182 (6) |

Table 2. Selected geometric parameters $\left(\AA^{\circ},^{\circ}\right)$

|  | $1.769(3)$ | S8-C10 | $1.812(3)$ |
| :--- | ---: | :--- | ---: |
| S1-C1 | $1.767(3)$ | $\mathrm{Cl}-\mathrm{C} 2$ | $1.359(2)$ |
| S1-C3 | $1.767(3)$ | $\mathrm{C} 3-\mathrm{C} 4$ | $1.356(2)$ |
| S2-C1 | $1.771(3)$ | $\mathrm{C} 5-\mathrm{C} 6$ | $1.360(2)$ |
| S2-C4 | $1.769(3)$ | $\mathrm{C} 7-\mathrm{C} 8$ | $1.525(3)$ |
| S3-C2 | $1.778(3)$ | $\mathrm{C} 9-\mathrm{C} 10$ | $1.523(2)$ |
| S3-C5 | $1.760(3)$ | D7A-C7 | $1.092(2)$ |
| S4-C2 | $1.769(3)$ | D7B-C7 | $1.092(2)$ |
| S4-C6 | $1.748(3)$ | D8A-C8 | $1.094(2)$ |
| S5-C3 | $1.848(4)$ | D8B-C8 | $1.101(2)$ |
| S5-C7 | $1.751(3)$ | D9A-C9 | $1.096(2)$ |
| S6-C4 | $1.832(4)$ | D9B-C9 | $1.100(2)$ |
| S6-C8 | $1.759(4)$ | D10A-C10 | $1.095(2)$ |
| S7-C5 | $1.817(4)$ | D10B-C10 | $1.102(2)$ |
| S7-C9 | $1.759(3)$ |  |  |
| S8-C6 | $93.8(2)$ | C4-C3-S1 | $116.9(2)$ |
| C1-S1-C3 | $93.3(2)$ | S5-C3-S1 | $118.2(2)$ |
| C1-S2-C4 | $94.0(1)$ | C3-C4-S6 | $122.0(2)$ |
| C2-S3-C5 | $94.6(2)$ | C3-C4-S2 | $117.6(1)$ |
| C2-S4-C6 | $103.4(2)$ | S6-C4-S2 | $120.0(2)$ |
| C3-S5-C7 | $96.4(2)$ | C6-C5-S7 | $129.6(2)$ |
| C4-S6-C8 | $101.1(2)$ | C6-C5-S3 | $116.9(2)$ |
| C5-S7-C9 | $100.1(2)$ | S7-C5-S3 | $113.2(2)$ |
| C6-S8-C10 | $122.1(2)$ | C5-C6-S8 | $127.2(2)$ |
| C2-C1-S1 | $123.1(2)$ | C5-C6-S4 | $116.7(2)$ |
| C2-C1-S2 | $114.6(2)$ | S8-C6-S4 | $115.7(2)$ |
| S1-Cl-S2 | $122.9(2)$ | C8-C7-S5 | $116.6(2)$ |
| C1-C2-S4 | $123.1(1)$ | C7-C8-S6 | $114.6(1)$ |
| C1-C2-S3 | $113.9(2)$ | C10-C9-S7 | $113.1(1)$ |
| S4-C2-S3 | $124.9(2)$ | C9-C10-S8 | $112.3(1)$ |
| C4-C3-S5 |  |  |  |

Table 3. $S \cdots S$ contacts $(\AA)$

| S1.. S2 ${ }^{\text {'1 }}$ | 3.593 (4) | $3.69 \dagger$ |
| :---: | :---: | :---: |
| S1.. S3 ${ }^{\text {/ii }}$ | 3.594 (4) | $3.62 \dagger$ |
| S4. . S7 ${ }^{\text {'iii }}$ | 3.393 (4) | $3.48 \dagger$ |
| S5.. $\mathrm{Sb}^{\text {'iv }}$ | 3.721 (4) | $3.69 \dagger$ |
| S7...S8 ${ }^{\text {fix }}$ | 3.503 (4) | $3.55 \dagger$ |

Symmetry codes: (i) $-x,-y,-z$; (ii) $-x-1,-y,-z$; (iii) $x+1, y, z$; (iv) $x-1, y, z$.
$\dagger$ Kobayashi et al. (1986).

Table 4. Refined scattering lengths ( $10^{-12} \mathrm{~cm}$ )

| C1 | $0.646(6)$ | S4 | $0.359(5)$ |
| :--- | :--- | :--- | :--- |
| C2 | $0.644(5)$ | S5 | $0.356(5)$ |
| C3 | $0.647(5)$ | S6 | $0.35(5)$ |
| C4 | $0.640(5)$ | S7 | $0.552(5)$ |
| C5 | $0.648(5)$ | S8 | $0.355(5)$ |
| C6 | $0.653(5)$ | D7A | $0.620(6)$ |
| C7 | $0.660(6)$ | D7B | $0.616(7)$ |
| C8 | $0.651(6)$ | D8A | $0.621(6)$ |
| C9 | $0.650(6)$ | D8B | $0.60(6)$ |
| C10 | $0.661(6)$ | D9A | $0.620(6)$ |
| S1 | $0.355(5)$ | D9B | $0.617(6)$ |
| S2 | 0.35445 | D10A | $0.624(6)$ |
| S3 | $0.346(5)$ | D10B | $0.626(6)$ |

A detailed description of the IPNS SCD has been published and described elsewhere (Schultz, 1993). Since the IPNS is a pulsed neutron source, data are obtained by the time-of-flight (TOF) Laue technique utilizing an area position-sensitive ${ }^{6} \mathrm{Li}-$ glass scintillator detector with an active area of $30 \times 30 \mathrm{~cm}^{2}$ and a wavelength range of $0.7-4.2 \AA$. Each detected neutron is stored in a three-dimensional histogram with coordinates $x$, $y, t$ corresponding to horizontal and vertical detector positions and the TOF, respectively. 32 histograms with different $\chi$ and $\varphi$ crystal settings were required to cover a unique quadrant of reciprocal space. Counting time per histogram (i.e. per $\chi$ and $\varphi$ setting) was approximately 6.5 h . Since all peaks were counted during each pulse no standards were required, only scaling between different crystal orientations.

The lattice parameters were refined from the observed positions and wavelengths of 1941 reflections. Bragg peaks were integrated in three dimensions ( $x, y, t$ ) about calculated peak positions from an orientation matrix. Data were corrected for the Lorentz factor, absorption, the incident source spectrum and the detector efficiency. Linear absorption coefficients at each wavelength were calculated from the sum of the wavelength-dependent true absorption and the wavelengthindependent total scattering: $\mu=0.355+0.0052 \times \lambda$, with $\mu$ in units of $\mathrm{cm}^{-1}$ and $\lambda$ in units of $\AA$.

The programs used to analyze the data were as follows: PEAKS, to search the histogram for Bragg peaks (Jacobson, 1986); BLIND, an autoindexing program for time-offlight data (Jacobson, 1986); LSQRS, to determine unit cell and orientation matrix least squares (Argonne program); $I N$ TEGRATE, which integrates Bragg peaks (Argonne program); $A N V R E D$, for data reduction (Argonne program); $A N V L S$, for least-squares structure refinement with multiwavelength data (Argonne program), based on ORFLS (Busing, Martin \& Levy, 1962); ORTEPIII, for molecular graphics (Burnett \& Johnson, 1996).

Supplementary data for this paper are available from the IUCr electronic archives (Reference: BR1170). Services for accessing these data are described at the back of the journal.

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## Two Related Benzylcamphor Dimers

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## Abstract

The crystal structures of two related dimeric 3-benzylcamphors, namely, $\left(1 R, 3 R, 4 R, 1^{\prime} R, 3^{\prime} R, 4^{\prime} R\right)-3,3^{\prime}-[(1 R, 2 R)$ -1,2-diphenylethylene]bis(1,7,7-trimethylbicyclo[2.2.1]-heptan-2-one), (1), and ( $1 R, 3 R, 4 R, 1^{\prime} R, 3^{\prime} S, 4^{\prime} R$ )-3, $3^{\prime}-$ [(1R,2S)-1,2-diphenylethylene]bis(1,7,7-trimethylbicyclo-[2.2.1]heptan-2-one), (2), both $\mathrm{C}_{34} \mathrm{H}_{42} \mathrm{O}_{2}$, are reported. Compound (1) possesses $C_{2}$ molecular symmetry; it adopts a conformation in the crystal in which the two phenyl rings are gauche with respect to one another and anti to the $\mathrm{C} 3-\mathrm{C} 4$ bond (terpene numbering) of the camphor ring system. Compound (2) lacks a molecular $C_{2}$ axis of symmetry; it adopts a conformation in the crystal where the two phenyl rings are anti with respect to one another, and in which one phenyl ring is partially in the endo cavity of the camphor ring system. In both compounds, C3 of both camphor ring systems (terpene numbering) is distorted to trigonal pyramidal from tetrahedral; the $\mathrm{C}-\mathrm{C}-\mathrm{C}$ bond angles about C 3 in both camphor rings are considerably expanded from idealized tetrahedral values and the dihedral angle between the benzyl C atom and the carbonyl O atom is reduced from the idealized value of $60^{\circ}$ to less than $48^{\circ}$.

## Comment

The ${ }^{13} \mathrm{C}$ NMR spectrum of the minor isomer (1) exhibits only ten resonances typical of $s p^{3}$-hybridized C atoms

